

L to H mode Transition: Parametric Dependencies of the Temperature Threshold

C. Bourdelle¹, L. Chôné^{1,2}, N. Fedorczak¹, X. Garbet¹, P. Beyer², J. Citrin^{1,3}, G. Dif-Pradalier¹, G. Fuhr², A. Loarte⁴, C.F. Maggi⁵, F. Millitello⁶, Y. Sarazin¹ and JET EFDA Contributors*

¹CEA, IRFM, F-13108 Saint Paul-lez-Durance, France. ²PIIM - UMR 7345 - Université d'Aix-Marseille - CNRS, 13397 Marseille Cedex 20, France ³FOM Institute DIFFER-Dutch Institute for Fundamental Energy Research, The Netherlands ⁴ITER Organization, Route de Vinon sur Verdon, 13115 St Paul Lez Durance, France ⁵Max Planck Institut für Plasmaphysik, EURATOM Association, Garching, Germany ⁶Culham Centre for Fusion Energy, Abingdon, UK JET-EFDA, Culham Science Centre, Abingdon, OX14 3DB, UK *See the Appendix of F. Romanelli et al., Proceedings of the 24th IAEA Fusion Energy Conference 2012, San Diego, US

Context

- The L to H mode transition occurs at **T threshold** which depends on **B, n and more** [Burrell 89, Suttrop 97, Hubbard 98, Righi 2000, Meakins 2010, etc]
- ExB is stabilizing** turbulence [Burrell review 97] more recently observed in AUG [P. Sauter NF2012], JET [Delabie this conf] and dithering transition identified as the result of an interplay between turbulence, zonal and mean E_r flows [Schmitz PRL 2012, Tynan NF 2013]
- The nature of the stabilized turbulence matters: Resistive Ballooning Modes key player in the highly collisional edge [Rogers-Drake-Zeiler 96-97]. Coherent with lower P_{th} observed for lower Z_{eff} in JET [Maggi NF2014, EPS2014, Bourdelle NF Letters 2014], AUG [Neu JNM2013], and earlier [Takizuka ITPA 2004]

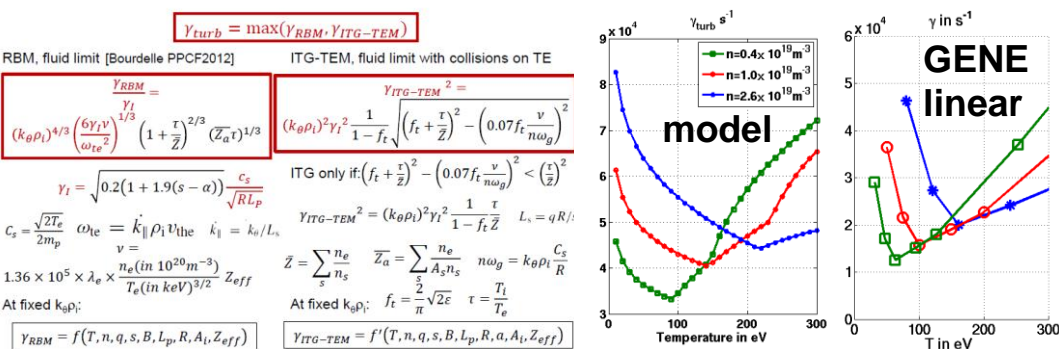
Driving idea

Transition when ExB time is shorter than turbulence time, or γ_{turb}/γ_E below a certain number. Account for both the ExB stabilization and the nature of the underlying turbulence.

Nature of Turbulence prior to Transition

- JET-ILW, prior to transition, $\rho=0.97$, linear stability analysis with GENE: **RBM unstable** [Bourdelle NF2014]

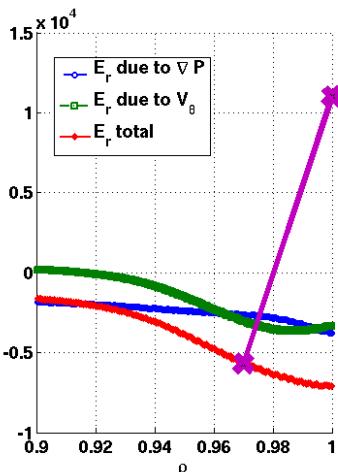
pulse	ρ	R/L_T	R/L_n	T	n	ν^*	q	s	Z_{eff}	B
82228	0.97	55	9	122	2.6	9.2	3.8	4.3	1.3	1.8



At low T RBM unstable, as T increases ITG-TEM take over. Min γ for T in experimental range. NB: α stabilization for $T \gg T_{exp}$

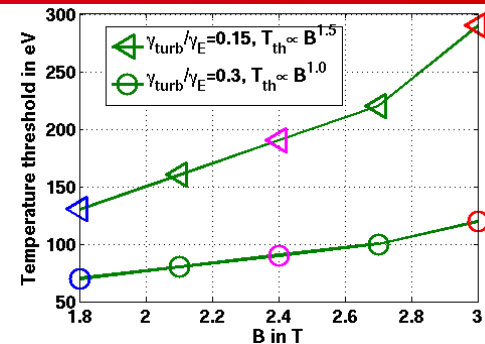
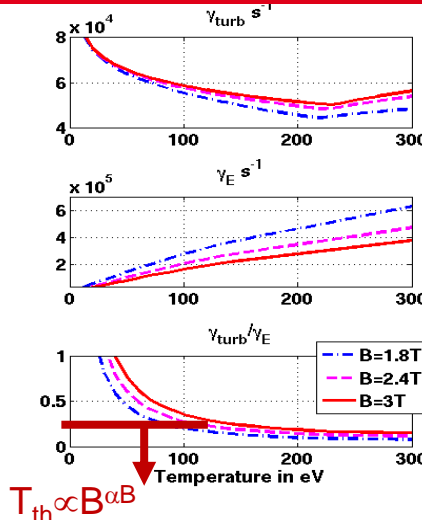
ExB Shear Derivation

- At $\rho=0.97$, E_r neo-classical: without V_ϕ (in JET at low NBI V_ϕ negligible, [Delabie this conf] in AUG no V_ϕ gradient in this region [Viezzer, NF13] with V_ϕ from banana to P-S regimes)
- At LCFS, E_r scales (at least) as $-3\nabla T_e$, assuming L_T constant across separatrix: $E_r(1) = 3 \frac{T(0.97)e^{(0.03a/L_T)}}{L_T}$



$$\gamma_E = \frac{\nabla E_r}{B} = \frac{E_r(0.97) - E_r(1)}{0.03 \times a \times B} = g(T, n, Z_{eff}, a, R, q, B, Z_i, L_T, L_n)$$

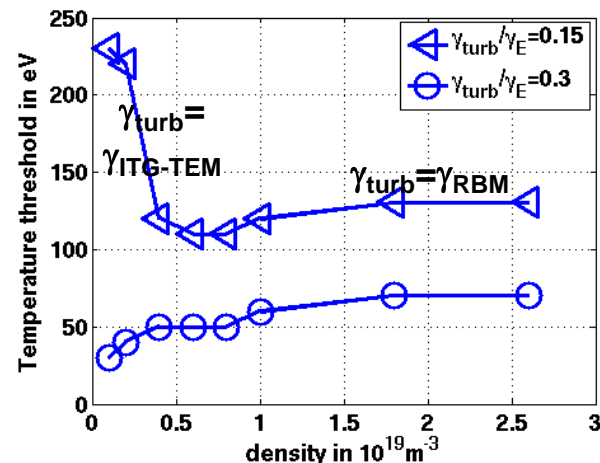
T threshold $\propto B$



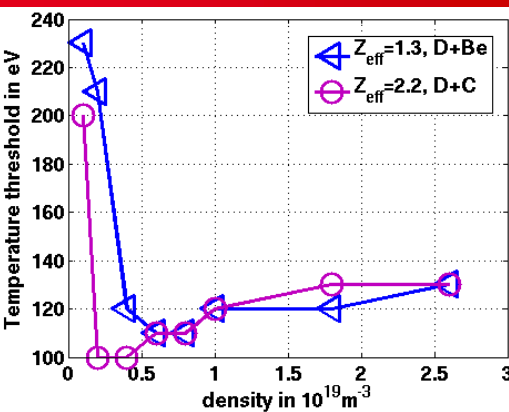
since: $P/S = n T V_{perp}$ [Ryter 96] and $T_{th} \propto B$ [Burrell 89], one gets: $P_{th} = n B S V_{perp}$ as given by [ITPA2004, 2008 and most scaling laws]

Minimum in density due to modified underlying turbulence

At low density, i.e. lower collisionalities, ITG-TEM dominate over RBM, $\gamma_{ITG-TEM}$ increases as n decreases
At high density, i.e. higher collisionalities, RBM dominate over ITG-TEM, γ_{RBM} increases as n increases

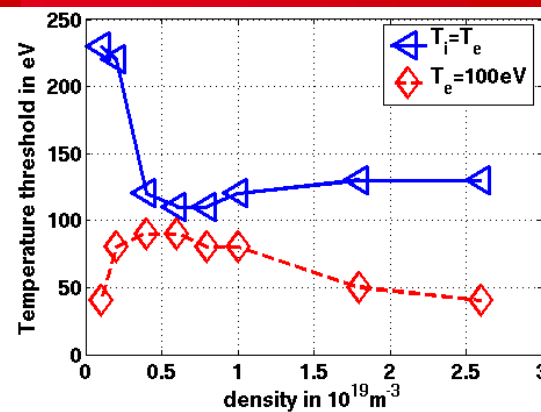


Lower Z_{eff} : min density shifted up



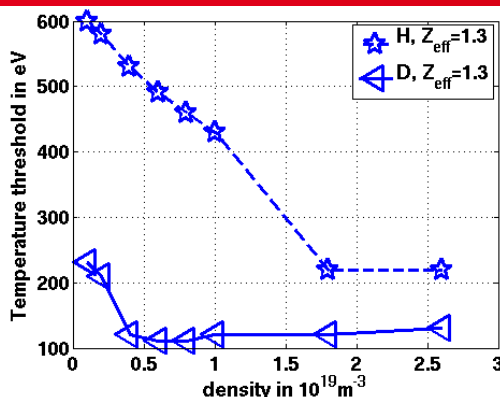
Coherent density min reappearing in JET-ILW with lower Z_{eff} [Maggi NF2014]

T_i increase at fixed T_e : min in density disappears

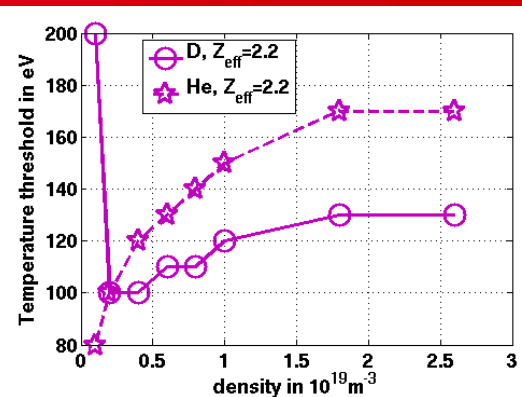


Coherent with AUG key role of ion heat flux at plasma edge [Ryter NF2014]

Higher threshold in H vs D and in He vs D, at fixed Z_{eff}



Coherent with isotope effect seen on P_{th} [Ryter96, Righi99]



Coherent with larger P_{th} in He vs D on DIII-D/Cmod [Gohil13]